

能量守恆法則

流體比熱為 c 、溫度為 T 時，領域內部熱能的變化率為

$$\int \frac{\partial}{\partial t} (\rho c T) d\Omega$$

從邊界流入熱能的流入率為

$$-\int \left(\rho c T v - k \frac{\partial T}{\partial n} \right) d\Gamma$$

兩者相等，得下列能量守恆法則

$$\int \frac{\partial}{\partial t} (\rho c T) d\Omega = -\int \left(\rho c T v - k \frac{\partial T}{\partial n} \right) d\Gamma$$

k 為熱傳導率。利用發散定理將上式右邊的邊界積分轉換成領域積分，可得

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$$\int \frac{\partial}{\partial t} (\rho c T) d\Omega = -\int \left(\frac{\partial \rho c T v_1}{\partial x_1} + \frac{\partial \rho c T v_2}{\partial x_2} \right) d\Omega + \int \left[\frac{\partial}{\partial x_1} \left(k \frac{\partial T}{\partial x_1} \right) + \frac{\partial}{\partial x_2} \left(k \frac{\partial T}{\partial x_2} \right) \right] d\Omega$$

考量質量守恆法則，得下列能量守恆方程式。

$$\rho \frac{\partial}{\partial t} (cT) + \rho v_1 \frac{\partial cT}{\partial x_1} + \rho v_2 \frac{\partial cT}{\partial x_2} - \left[\frac{\partial}{\partial x_1} \left(k \frac{\partial T}{\partial x_1} \right) + \frac{\partial}{\partial x_2} \left(k \frac{\partial T}{\partial x_2} \right) \right] = 0$$

當比熱 c 與熱傳導率 k 一定時，得

$$\frac{\partial T}{\partial t} + v_1 \frac{\partial T}{\partial x_1} + v_2 \frac{\partial T}{\partial x_2} - K \left(\frac{\partial^2 T}{\partial x_1^2} + \frac{\partial^2 T}{\partial x_2^2} \right) = 0$$

溫度擴散率 $K = k / (\rho c)$ 。

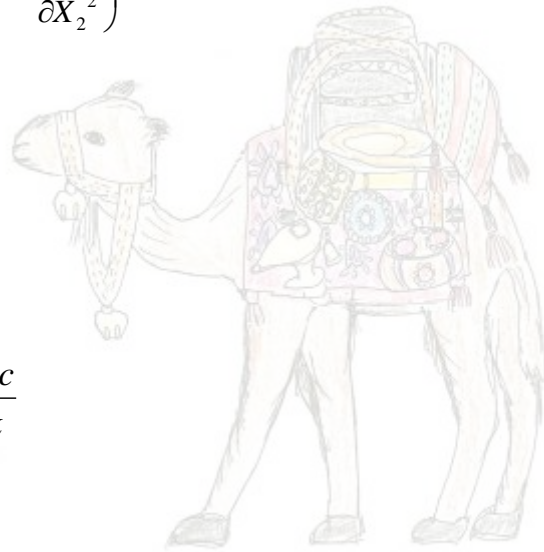
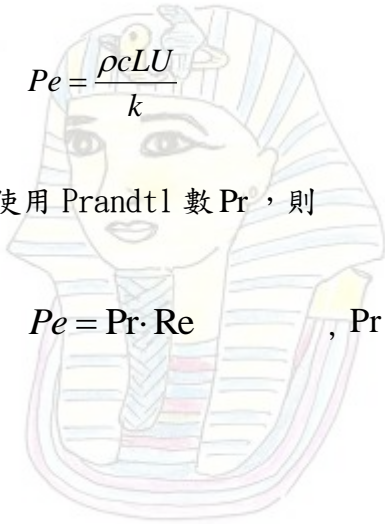
無因次化能量守恆方程式為

$$\frac{\partial \theta}{\partial \tau} + u_1 \frac{\partial \theta}{\partial X_1} + u_2 \frac{\partial \theta}{\partial X_2} + \frac{1}{Pe} \left(\frac{\partial^2 \theta}{\partial X_1^2} + \frac{\partial^2 \theta}{\partial X_2^2} \right) = 0$$

$$Pe = \frac{\rho c L U}{k}$$

若使用 Prandtl 數 Pr ，則

$$Pe = Pr \cdot Re, \quad Pr = \frac{\mu c}{k}$$



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回應用邊界元素法解析海洋擴散

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